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# ENC-2020-0426 NUMERICAL SOLUTION OF INTERNAL FLOW THROUGH A DE LAVAL NOZZLE BASED ON THE EULER MODEL WITH THE SU2 CODE: VERIFICATION AND VALIDATION

## Giovanne Deni Iorio

Postgraduate Program in Mechanical Engineering (PGMEC), Federal University of Paraná (UFPR), Curitiba, PR, Brazil giovanne.iorio@ufpr.br

## **Guilherme Bertoldo**

Department of Physics, Statistics and Mathematics (DAFEM), Federal University of Technology – Paraná (UTFPR), Campus Francisco Beltrão, ZIP Code 85601-970, Francisco Beltrão, PR, Brazil gbertoldo@utfpr.edu.br

## Carlos Henrique Marchi

Laboratory of Numerical Experimentation (LENA), Department of Mechanical Engineering (DEMEC), Federal University of Paraná (UFPR), P.O. Box 19040, ZIP Code 81531-980, Curitiba, PR, Brazil chmcfd@gmail.com

Abstract. The present work performs the verification of the numerical solution of the internal flow through a de Laval nozzle obtained with the SU2 code and the validation of Euler's mathematical model. The flow was solved numerically in several meshes, where the finest one had 2048x1024 volumes. The discharge coefficient is the variable of interest, a coefficient that characterizes the flow and pressure losses. Comparisons of the numerical discharge coefficient with analytical and experimental results showed relative differences of 0.0257% and 0.379%, respectively, for the most refined simulated mesh. Comparisons also were made with the numerical solution of another computational code, finding relative differences of 0.000185%. Estimates of the numerical error were calculated based on error estimators. According to the GCI estimator, the numerical uncertainty does not exceed 0.001% of the numerical solution in the finest simulated mesh. Extrapolated solutions with Repeated Richardson Extrapolation and the convergent solution are also presented.

Keywords: nozzle, verification, validation, CFD, SU2

## 1. INTRODUCTION

Verification and validation (V&V) are procedures that make it possible to assess the accuracy and reliability of a numerical solution. The first is the process that quantifies the numerical error and the second is the process that quantifies the modeling error caused by the limitations of the mathematical models to represent the real phenomenon (Roache, 2009).

This work aims to verify the solution of Euler's numerical model implemented in the SU2 code (Economon et al., 2016) and to validate Euler's mathematical model in the solution of fluid flow through a de Laval nozzle. A conical nozzle with angles of 45° in the convergent section and 15° in the divergent section will be considered. This nozzle has experimental data published by Back et al. (1965). The variable of interest in this study will be the discharge coefficient, a coefficient that characterizes the flow and pressure losses.

According to the ASME V&V 20-2009 standard, the verification of the solution must be preceded by code verification. That is why Van der Weide and Economon (2019) applied the method of manufactured solutions to perform the code verification of SU2 (version 7.0.0) and concluded that the code is error-free. The process of verifying numerical solutions and validating models obtained/implemented by the SU2 code was addressed by other authors (Palacios et al., 2013; Economon et al., 2016; Gori et al., 2017; Becker and Granzoto, 2018; Mishra et al., 2019; Castro, 2019). However, V&V of the SU2 code is usually performed for the Navier-Stokes mathematical model or turbulence models and the errors or numerical error estimates of the solutions obtained are generally not presented. Of the studies cited, only in the works of Mishra et al. (2019) and Castro (2019) the numerical error or its estimate were presented. Even though the SU2 code shows certain accuracy for a particular application, it does not mean that the same accuracy will be achieved in another application. Thus, a numerical simulation code must be evaluated in the largest possible number of applications.

In the present work, the estimated numerical uncertainty for the variable of interest will be calculated using the GCI (Grid Convergence Index) estimator (Roache, 1994), using the convergent estimator (Marchi and Silva, 2002) and with the estimator based on the Repeated Richardson Extrapolation (RRE) (Martins, 2013). In addition, the numerical solutions will be extrapolated with RRE and with the convergent estimator (Marchi et al., 2013).

In the following section, the flow simulation method and a definition of the variable of interest are presented. The procedures used for mesh generation and verification and validation are also shown. Section 3 presents the main results. Finally, a conclusion to this work is presented.

## 2. METHODOLOGY

## 2.1 Flow simulation

The flow is considered inviscid, non-reactive, axisymmetric and modeled by Euler's equations, which can be represented by (Hirsch, 2007)

$$\frac{\partial \mathbf{U}}{\partial t} + \nabla \cdot \mathbf{F}^c(\mathbf{U}) = 0 \tag{1}$$

where U is the vector of conservative variables and  $\mathbf{F}^c$  is the advective flux, given by

$$\mathbf{U} = \begin{cases} \rho \mathbf{v} \\ \rho \mathbf{v} \\ \rho E \end{cases}, \ \mathbf{F}^c = \begin{cases} \rho \mathbf{v} \\ \rho \mathbf{v} \otimes \mathbf{v} + \mathbf{\bar{I}} p \\ \rho E \mathbf{v} + p \mathbf{v} \end{cases}$$
 (2)

In Eq. (1) t is time and in Eq. (2)  $\rho$  is the density,  $\mathbf{v}$  is the velocity vector, E is the total energy per unit mass and p is the pressure. Furthermore,  $\mathbf{I}$  is a 2x2 identity matrix and  $\otimes$  represents the tensor product. In Euler's system of equations, the equations represent the conservation of mass, the conservation of momentum and the conservation of total energy. Euler's equations have as dependent variables  $\rho$ ,  $\mathbf{v}$ , p and E; one can close the system of equations by assuming a constitutive relationship or state equation for the fluid, for example

$$p = (\gamma - I)\rho[E - 0.5(\mathbf{v} \cdot \mathbf{v})] \tag{3}$$

where  $\gamma$  is the ratio of specific heats. In the SU2 code, the discretization of the Euler's equation is performed through the Finite Volume Method in a vertex-based mesh (Economon et al., 2016). The advective fluxes are evaluated at the midpoint of the edges.

The discharge coefficient  $(C_d)$ , which is the variable of interest in this work, characterizes the flow and pressure losses and it is defined as the ratio between the actual mass flow rate  $(\dot{m})$  and the ideal mass flow rate  $(\dot{m}_{ideal})$  (Anderson Junior, 2003) of the combustion gases. The equations of the ideal mass flow rate and the actual mass flow rate are shown in Eq. (4) and Eq. (5), respectively.

$$\dot{m}_{ideal} = p_0 A_t \sqrt{\frac{\gamma}{RT_0} \left(\frac{2}{\gamma + I}\right)^{\frac{\gamma + I}{\gamma - I}}} \tag{4}$$

$$\dot{m} = \int_{A} \rho u dA \tag{5}$$

In Eq. (4),  $\gamma$  is the ratio of specific heats, R is the gas constant and  $A_t$  is the throat area.  $T_0$  and  $p_0$  are the temperature and pressure of stagnation. In Eq. (5), u is the axial component of the velocity vector,  $\rho$  is the density and A is the nozzle exit area. The integral present in Eq. (5) was numerically approximated using the Trapezoidal Rule.

In the simulations, the physical and geometric parameters of the Back et al. (1965) experiment with the nozzle BGM45-15 were considered, whose profile is illustrated in Fig. 1. Table 1 presents the flow parameters.

Table 1. Flow parameters.

| Parameter                 | Value                     |
|---------------------------|---------------------------|
| Temperature of stagnation | 833.333 K (1500 °R)       |
| Pressure of stagnation    | 1725.068 kPa (250.2 psia) |
| External pressure         | 0 Pa                      |
| Ratio of specific heats   | 1.4                       |
| Gas constant              | 287.058 J/kg.K            |

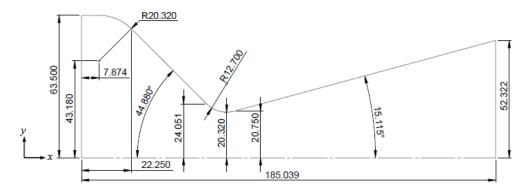


Figure 1. BGM45-15 nozzle profile with dimensions in millimeters.

The boundary conditions on the nozzle's wall is of adiabacity and impermeability, e.g.

$$\mathbf{n} \cdot \mathbf{v} = 0, \ \frac{\partial p}{\partial n} = 0, \ \frac{\partial T}{\partial n} = 0$$
 (6)

where T is the temperature and  $\mathbf{n}$  is the normal vector. The first equation represents the fluid slipping condition, since the inviscid flow will be considered.

On the symmetry line, the radial velocity v is zero and adiabaticity and impermeability conditions are also considered for this boundary.

On the inlet of the nozzle, it is assumed that

$$\frac{\partial^2 u}{\partial v^2} = 0, \quad v = 0 \tag{7}$$

where x is the axial direction. It is also considered that the inlet temperature and the inlet pressure are functions of the stagnation properties  $T_0$  and  $p_0$  presented in Tab. 1. These functions are given by (Sutton and Biblarz, 2016)

$$T = T_0 [I + \frac{1}{2} (\gamma - I) M^2]^{-1}$$
 (8)

$$p = p_0 \left[ 1 + \frac{1}{2} (\gamma - I) M^2 \right]^{\frac{-\gamma}{\gamma - I}}$$
 (9)

where M is the Mach number and  $\gamma$  is the ratio of specific heats.

On the nozzle's outlet boundary, no boundary conditions are necessary, as the flow is supersonic, a situation in which the flow depends only on the upstream characteristics of the flow. However, numerically, a boundary condition is required, and the condition used is the extrapolation of the properties of the interior of the domain to the exit boundary.

For the simulations, the version 7.0.6 of the SU2 code with multiprocessing support for the Linux operating system was used. The simulations were performed in a computer with a processor Intel Core i5-9600K of 4.60 GHz and 16 GB of RAM memory.

# 2.2 Mesh generation

Non-orthogonal structured meshes were used in the simulations, which were generated using constant and simultaneous refinement ratio in the radial and axial directions. For the mesh generation, the code GMSH version 4.5.6 (Geuzaine and Remacle, 2009) for Linux was used.

The base mesh used in the simulations has 32 control volumes in the axial direction and 16 volumes in the radial direction. Figure 2 shows this base mesh. From this mesh, using a refinement ratio of two, six other meshes were generated. Thus, the following meshes were created: 32x16, 64x32, 128x64, 256x128, 512x256, 1024x512 and 2048x1024 volumes.

# 2.3 Errors and estimated errors

Numerical errors are always present in the flow solutions obtained through computational fluid dynamics (CFD). The main sources of these errors are: truncation errors  $E_h$ , round-off errors  $E_{\pi}$  and iteration errors  $E_i$  (Roache, 2009). In general, the numerical error  $E_n$ , present in the numerical solution of a generic variable of interest  $\phi$ , can be symbolically represented by

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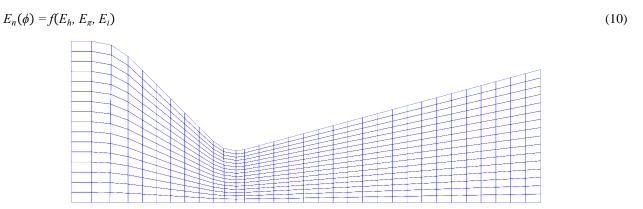


Figure 2. Base mesh with 32x16 volumes.

The numerical error can also be defined as the difference between the numerical solution  $\phi$  and the exact analytical solution  $\Phi$  of a variable of interest (Ferziger et al., 2020), i.e.,

$$E_n(\phi) = \Phi - \phi \tag{11}$$

In this work, care was taken to minimize round-off and iteration errors so that these errors could be considered negligible. Round-off errors were minimized by using double precision variables. Iteration errors were minimized by defining as a stop criterion of the iterative process as twice the number of iterations required for the residue of the mass conservation equation to be less than or equal to the tolerance of 10<sup>-14</sup>. Hence, the numerical error of the numerical solution is represented only by the truncation error, that is,

$$E_h(\phi) = E_n(\phi) = \Phi - \phi \tag{12}$$

In this condition, the numerical error is denominated discretization error. Note that in practice,  $\Phi$  is not known, so the numerical error must be estimated  $(U_n)$ . Thereby,

$$U_n = \phi_{\infty} - \phi \tag{13}$$

where  $\phi_{\infty}$  is an estimate of the exact solution.

The discretization error can be effectively estimated after obtaining the numerical solution. Some of the discretization error/uncertainty estimators are the GCI estimator (Roache, 1994), convergent estimator (Marchi and Silva, 2002) and RRE-based estimator (Martins, 2013). These estimators are used in this work.

In the verification process, the ASME V&V 20-2009 standard recommends performing the apparent order convergence analysis. As the mesh is refined, the apparent order  $p_U$  is expected to tend to the asymptotic order  $p_0$  of the discretization error of the equations that model the problem. Divergences may indicate that other error sources are being generated by the code, inappropriate boundary conditions or incorrect initial conditions. In this study, it is expected that  $p_U$  converges to two.  $p_U$  is calculated with Eq. (14), where  $\phi_{SC}$  is the solution obtained in a supercoarse mesh,  $\phi_C$  is the solution obtained in a coarse mesh and  $\phi_E$  is the solution obtained in a fine mesh and r is the refinement ratio.

$$p_U = \frac{log\left(\left|\frac{\phi_C - \phi_{SC}}{\phi_F - \phi_C}\right|\right)}{log(r)} \tag{14}$$

According to the ASME V&V 20-2009 standard, in addition to the numerical error  $E_n$ , a numerical solution  $\phi$  is also subject to the modeling errors  $E_{model}$  and errors due to the uncertainty of the input data of the simulation ( $E_{input}$ ). The modeling error is expected to be contained in the range given by

$$(E \pm U_{val}) \tag{15}$$

in other words,

$$E_{model} \in [E - U_{val}, E + U_{val}] \tag{16}$$

In Eq. (15) and Eq. (16), the symbols E and  $U_{val}$  are defined by the ASME V&V 20-2009 standard as validation metrics. The former is denominated comparison error and it is given by the difference between the numerical solution  $\phi$  and the experimental result X.  $U_{val}$ , the validation standard uncertainty, can be calculated as

$$U_{val} = \sqrt{U_{num}^2 + U_{input}^2 + U_{exp}^2}$$
 (17)

where  $U_{num}$ ,  $U_{input}$  and  $U_{exp}$  are estimates of the standard numerical uncertainty, the standard uncertainty of the simulation input data and the standard uncertainty of the experimental error, respectively. Equation (17) is valid for cases where uncertainties are independent. The methodology for calculating  $U_{input}$  is presented in the ASME V&V 20-2009 standard. In addition, this standard relates  $U_{num}$  to the numerical error estimate of the GCI estimator ( $U_{GCI}$ ) as follows

$$U_{num} = \frac{U_{GCI}}{k}, \quad 1.1 \le k \le 1.15$$
 (18)

Back et al. (1965) presented the experimental result (X) for the discharge coefficient  $(C_d)$ . However, the experimental results were not tabulated in that study, being presented only in graphic form. Therefore, the code WebPlotDigitizer 4.2 (Marin et al., 2017) was used to extract the results from the graph. Hence, the value of the discharge coefficient obtained is

$$C_d = 0.9777 \pm 0.0056 \tag{19}$$

The uncertainty presented in the value of the experimental discharge coefficient encompasses both the experimental uncertainties and the uncertainty of reading the values in the graph.

Kliegel and Levine's (1969) analytical solution for the discharge coefficient, which assumes the flow as irrotational and isentropic, will be used for comparison with the values of the discharge coefficient obtained from the numerical simulations. For the problem treated in this work, the value obtained for the discharge coefficient by the method of Kliegel and Levine (1969) is

$$C_d = 0.981653876 \tag{20}$$

This value was calculated using quadruple precision.

The numerical solutions obtained in this present work will also be compared with the numerical discharge coefficient by Araki and Marchi (2017). Araki and Marchi (2017) studied the inviscid flow through the nozzle BGM45-15 and obtained the solution

$$C_d = 0.98140 \pm 0.00002 \tag{21}$$

for the ratio of specific heats 1.4.

# 3. RESULTS AND DISCUSSION

Table 2 presents the main results of the simulations. In this table,  $N_x$  is the number of volumes in the axial direction and  $N_y$  is the number of volumes in the radial direction. Moreover, "Iterations" indicates the number of iterations performed, "Time" represents the total simulation time, "RAM" means the maximum RAM memory consumption of each simulation,  $C_d$  is the discharge coefficient and  $p_U$  is the apparent order, calculated with Eq. (14). From this table, it is observed that the apparent order converges to two as the mesh is refined. Figure 3 shows the Mach number field obtained in the mesh  $m_7$ , which had a total of 2097152 volumes.

Table 2. General characteristics of the simulations.

| Mesh  | $N_x$ | $N_{y}$ | Iterations | Time      | RAM (MB) | $C_d$          | $p_U$ |
|-------|-------|---------|------------|-----------|----------|----------------|-------|
| $m_I$ | 32    | 16      | 1436       | 1.49 s    | 113.92   | 1.01206217E+00 | -     |
| $m_2$ | 64    | 32      | 1844       | 5.56 s    | 120.19   | 9.88127628E-01 | -     |
| $m_3$ | 128   | 64      | 2538       | 30.73 s   | 148.87   | 9.83761632E-01 | 2.45  |
| $m_4$ | 256   | 128     | 3398       | 4.20 min  | 249.24   | 9.82109643E-01 | 1.40  |
| $m_5$ | 512   | 256     | 5354       | 31.06 min | 635.93   | 9.81543049E-01 | 1.54  |
| $m_6$ | 1024  | 512     | 10428      | 4.81 h    | 2176.48  | 9.81427137E-01 | 2.29  |
| $m_7$ | 2048  | 1024    | 24535      | 1.91 day  | 8297.22  | 9.81401820E-01 | 2.19  |



Figure 3. Mach number field for the mesh with 2048x1024 volumes.

The discharge coefficient is presented with nine significant figures, as this is the maximum number of figures that was required in the final representation of the solutions with error estimation. In Roy and Oberkampf's (2011) framework for verification, validation and uncertainty quantification, two significant figures are used to report the error estimate and the solution is reported with the corresponding number of decimal places.

Table 3 shows comparisons between the numerical discharge coefficient and the analytical solution calculated by the method of Kliegel and Levine (1969). It is noticed that the absolute differences shown in Tab. 3 reduce with the mesh refinement, but from mesh  $m_6$ , the difference increases. This may have occurred because Kliegel and Levine's (1969) analytical solution does not necessarily represent the analytical solution of Euler's mathematical model since in its derivation the flow was considered irrotational and isentropic.

Table 3. Comparisons between numerical solutions and the irrotational and isentropic analytical solution of the discharge coefficient.

| Case       | $C_d$          | Difference | Relative difference |
|------------|----------------|------------|---------------------|
| $m_1$      | 1.01206217E+00 | 3.04E-02   | 3.10%               |
| $m_2$      | 9.88127628E-01 | 6.47E-03   | 0.659%              |
| $m_3$      | 9.83761632E-01 | 2.11E-03   | 0.215%              |
| $m_4$      | 9.82109643E-01 | 4.56E-04   | 0.0464%             |
| $m_5$      | 9.81543049E-01 | 1.11E-04   | 0.0113%             |
| $m_6$      | 9.81427137E-01 | 2.27E-04   | 0.0231%             |
| $m_7$      | 9.81401820E-01 | 2.52E-04   | 0.0257%             |
| Convergent | 9.81394063E-01 | 2.598E-04  | 0.02647%            |
| RRE        | 9.81393831E-01 | 2.600E-04  | 0.02649%            |

The estimated numerical uncertainty/error of the numerical solution in the finest mesh obtained with the estimators GCI, convergent and RRE-based estimator, as well as the solutions extrapolated with the convergent solution and with RRE are presented in Tab. 4. Table 4 also shows the expression of the final solutions with the estimated numerical error. One can notice that the estimated error of the convergent solution and the solution extrapolated with RRE are, respectively, two and three orders of magnitude smaller than the estimated error obtained with the GCI estimator for the mesh  $m_7$ .

Table 4. Numerical discharge coefficient and its error/uncertainty estimation.

| Method     | $C_d$          | Estimated error | Representation              |
|------------|----------------|-----------------|-----------------------------|
| GCI        | 9.81401820E-01 | 1.05E-05        | $0.981402 \pm 0.000010$     |
| Convergent | 9.81394063E-01 | 6.82E-07        | $0.98139406 \pm 0.00000068$ |
| RRE        | 9.81393831E-01 | 4.82E-08        | 0.981393831 + 0.000000048   |

The solution obtained with the GCI method is close to that obtained by Araki and Marchi (2017) for the ratio of specific heats 1.4. Araki and Marchi (2017) obtained the solution  $0.98140 \pm 0.00002$ , also using a second order scheme and the Euler's mathematical model. They obtained this solution using a mesh with only 720x80 volumes; however, their numerical uncertainty is higher. Table 5 presents the comparisons between the numerical discharge coefficient obtained in this work and the Araki and Marchi's (2017) numerical solution. It can be noticed that the Araki and Marchi's (2017) solution range includes all the numerical solutions presented in Tab. 4. ASME V&V 20-2009 standard recommends comparing the numerical solution with ones obtained by other codes for the solution verification process. These kind of comparisons are denominated code-to-code comparison.

| Case       | $C_d$          | Difference | Relative difference |
|------------|----------------|------------|---------------------|
| $m_I$      | 1.01206217E+00 | 3.07E-02   | 3.12%               |
| $m_2$      | 9.88127628E-01 | 6.73E-03   | 0.686%              |
| $m_3$      | 9.83761632E-01 | 2.36E-03   | 0.241%              |
| $m_4$      | 9.82109643E-01 | 7.10E-04   | 0.0723%             |
| $m_5$      | 9.81543049E-01 | 1.43E-04   | 0.0146%             |
| $m_6$      | 9.81427137E-01 | 2.71E-05   | 0.00277%            |
| $m_7$      | 9.81401820E-01 | 1.82E-06   | 0.000185%           |
| Convergent | 9.81394063E-01 | 5.94E-06   | 0.000605%           |
| RRE        | 9.81393831E-01 | 6.17E-06   | 0.000629%           |

Table 5. Comparisons between numerical solutions and the Araki and Marchi's numerical solution.

The comparisons between the numerical discharge coefficient and the experimental result obtained by Back et al. (1965) are presented in Tab. 6. In this case, the difference also reduces with the refinement of the mesh. The RRE solution has a smaller difference than the other ones.

Table 6. Comparisons between numerical solutions and the experimental result of the discharge coefficient.

| Case           | $C_d$          | Comparison error | Relative difference |
|----------------|----------------|------------------|---------------------|
| $m_1$          | 1.01206217E+00 | 3.44E-02         | 3.51%               |
| $m_2$          | 9.88127628E-01 | 1.04E-02         | 1.07%               |
| $m_3$          | 9.83761632E-01 | 6.06E-03         | 0.620%              |
| $m_4$          | 9.82109643E-01 | 4.41E-03         | 0.451%              |
| $m_5$          | 9.81543049E-01 | 3.84E-03         | 0.393%              |
| $m_6$          | 9.81427137E-01 | 3.73E-03         | 0.381%              |
| $\underline{}$ | 9.81401820E-01 | 3.70E-03         | 0.379%              |
| Convergent     | 9.81394063E-01 | 3.6941E-03       | 0.37783%            |
| RRE            | 9.81393831E-01 | 3.6938E-03       | 0.37781%            |

From Table 6, the validation metric E for the mesh  $m_7$  is 0.0037. The validation standard uncertainty  $U_{val}$  for the most refined simulated mesh was calculated as

$$U_{val} = \sqrt{\left(\frac{U_{GCI}}{k}\right)^2 + U_{input}^2 + U_{exp}^2} = 0.0056$$
 (22)

where k equal to 1.1 was used and the  $U_{GCI}$  estimated numerical error for the mesh  $m_7$  is shown in Tab. 4.  $U_{input}$  was calculated according to the methodology presented in ASME V&V 20-2009. The effects of temperature and pressure of stagnation were considered, and it resulted in  $U_{input} = 5.00$ E-16, indicating a low sensitivity of the discharge coefficient about these input parameters. It is noted that the predominant uncertainty in the validation standard uncertainty is the experimental uncertainty. As  $|E| < U_{val}$ , the combination of numerical, experimental and input data uncertainties is at the same order as the modeling error, i.e., the modeling error is within the noise level of the uncertainties, not allowing to evaluate whether the difference between the numerical solution and the experimental result is caused by the modeling error or other sources of error.

Figure 4 shows the pressure ratio distribution on the nozzle wall and at the symmetry line of the numerical solution of the mesh  $m_7$ . This figure also presents the pressure ratio distribution from the experimental wall pressures results of Back et al. (1965), which were presented only in graphic form by the authors, but whose values were extracted and tabulated by Radtke et al. (2013). As noted in Fig. 4, the numerical results of the wall pressure ratio are qualitatively similar to the experimental results, where the most significant differences are in the throat area.

The Mach number distribution on the nozzle wall and at the symmetry line are shown in Fig. 5. One can see an abrupt change in the Mach number in the symmetry line at position x = 0.1476. This is the position in which the shock wave meets the symmetry line, as seen in Fig. 3. The shock wave formed in this nozzle has experimental evidence given by Back and Cuffel (1966). According to the graph of pressure ratio distribution at the symmetry line presented by Back and Cuffel (1966), the position that the shock wave reaches the line of symmetry was calculated as  $x = 0.1461 \pm 0.0017$ , where the uncertainty considers the experimental pressure uncertainty and the uncertainty of reading the values in the graph. The abrupt change also happens for the other flow properties, as can be seen in Fig. 4 for the pressure ratio at the symmetry line and in Fig. 6 for the pressure field obtained in the mesh  $m_7$ . The temperature and density distributions are not shown in this work, as they have the same behavior as the pressure ratio.

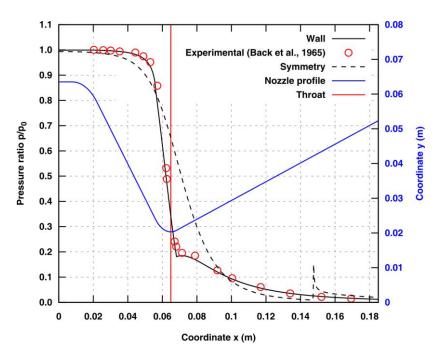


Figure 4. Comparison between the experimental and numerical results of the pressure ratio.

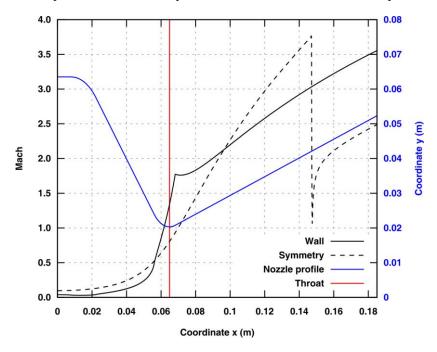


Figure 5. Numerical results of the Mach number on the nozzle wall and at the symmetry line.



Figure 6. Pressure field for the mesh with 2048x1024 volumes.

## 4. CONCLUSION

Comparisons between numerical discharge coefficients and analytical and experimental results were presented in this work and the differences showed reductions with mesh refinement. The relative differences for the most refined mesh simulated were 0.0257% in relation to the analytical result and 0.379% to the experimental result. The numerical solutions of the discharge coefficient also were compared to the numerical solution obtained by another computational code using the same mathematical model. In this case, the relative difference for the finest simulated mesh was 0.000185%. There was also a convergence of the apparent order to the asymptotic order, which is one of the code verification requirements proposed by the ASME V&V 20-2009 standard. It is estimated that the exact numerical solution for the discharge coefficient is contained in the range given by  $0.981402 \pm 0.000010$ . As for the modeling error, it is expected to be contained in the interval given by  $0.38\% \pm 0.57\%$  of the experimental result, where the first term is the comparison error and the second is the validation standard uncertainty, both divided by the experimental result of the discharge coefficient. In this case, the validation standard uncertainty is greater than the comparison error, indicating that the modeling error is within the noise level of the experimental, input and numerical uncertainties.

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