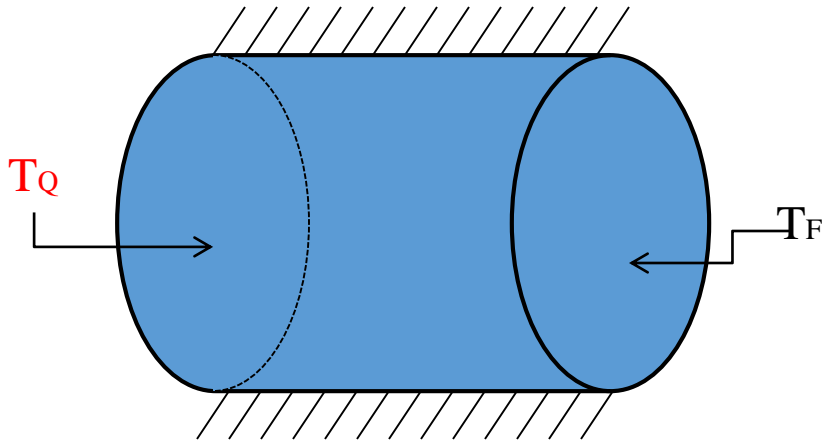


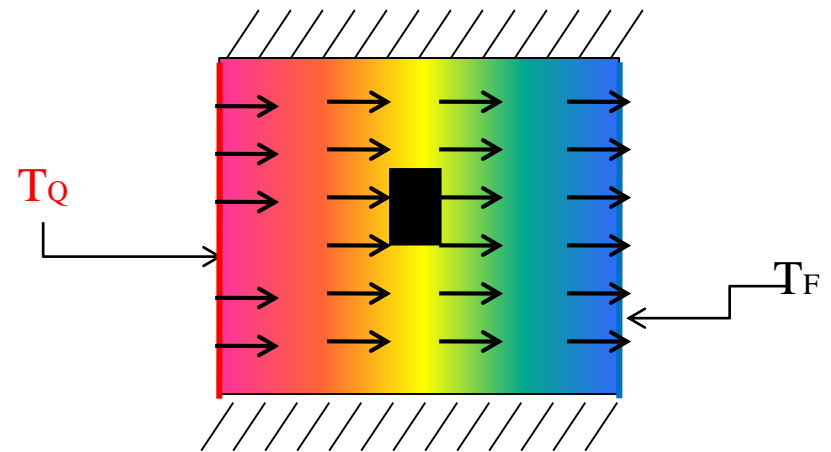
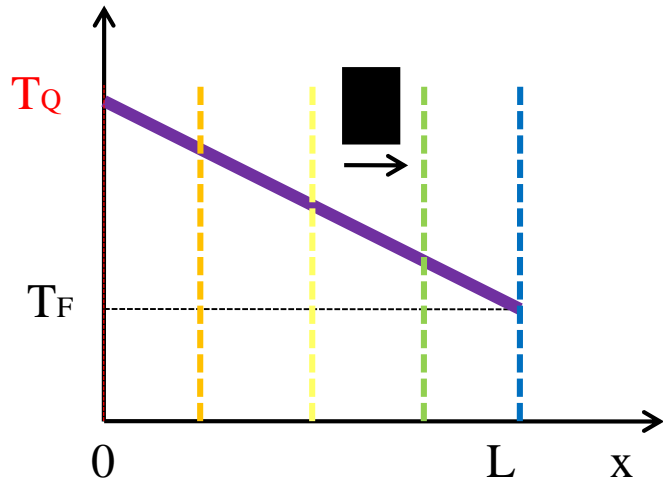
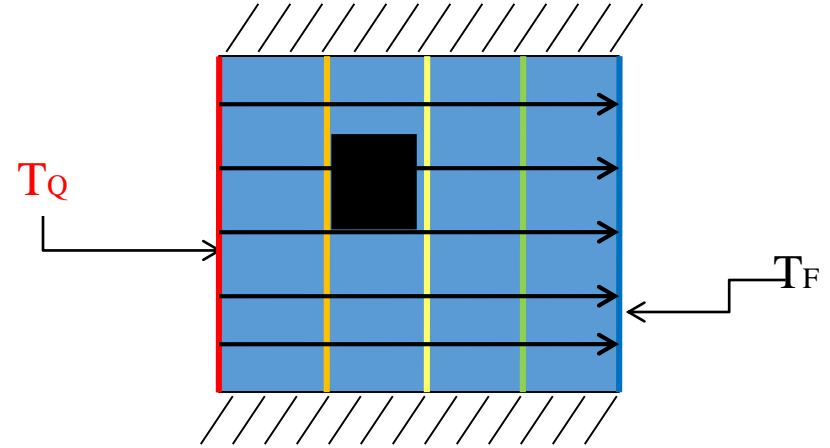
Condução Bidimensional em Regime Permanente

Recapitulação: Isotermas e Isofluxos – 1D

isolado



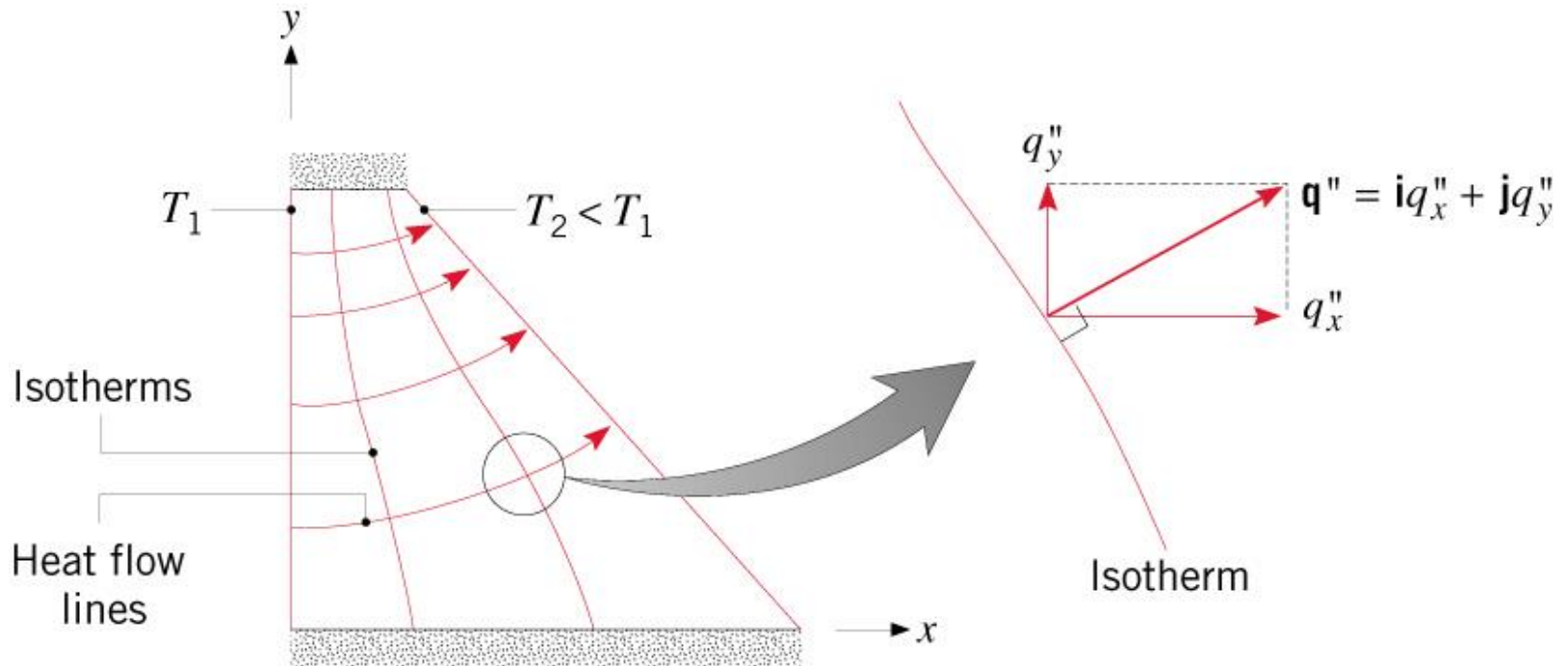
isolado



Introdução

- - $T(x,y)$
 - q_x'' , q_y''

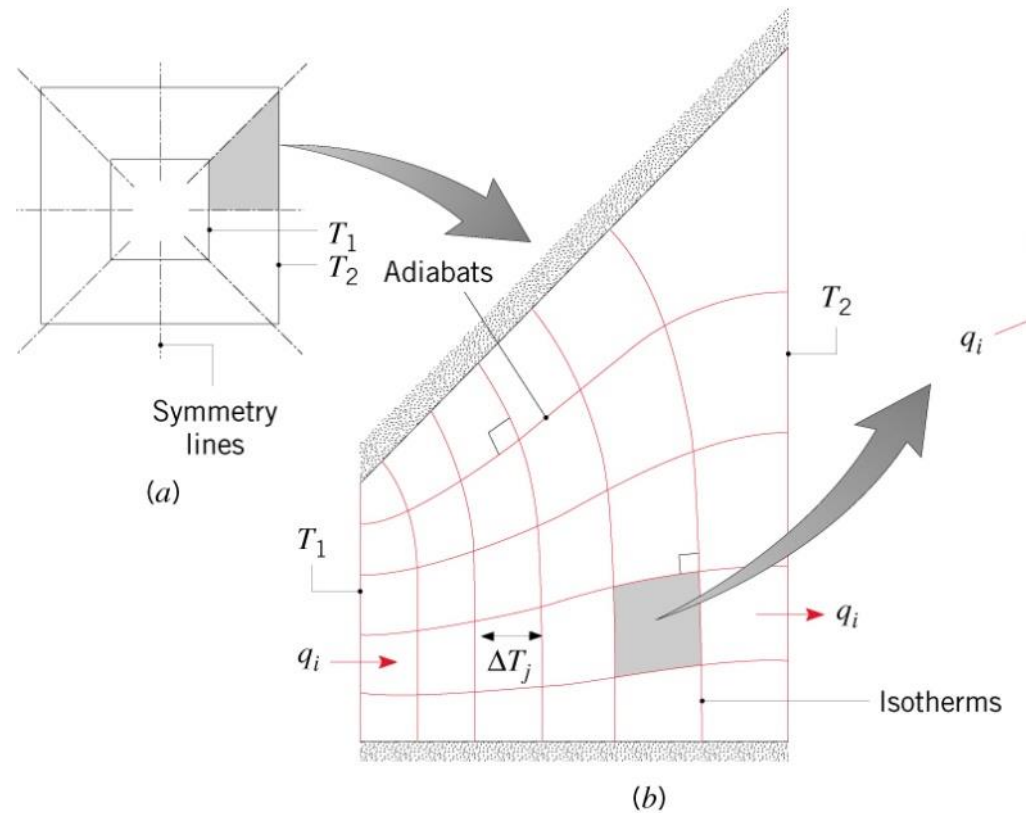
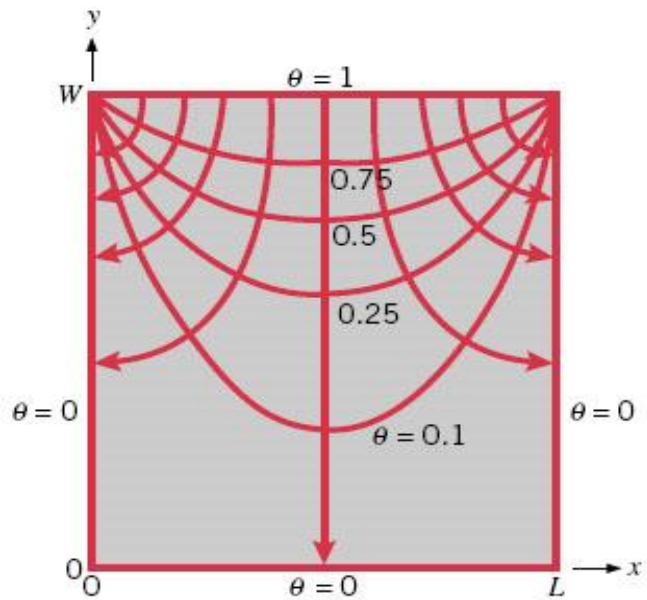
•Isotermas e isofluxos



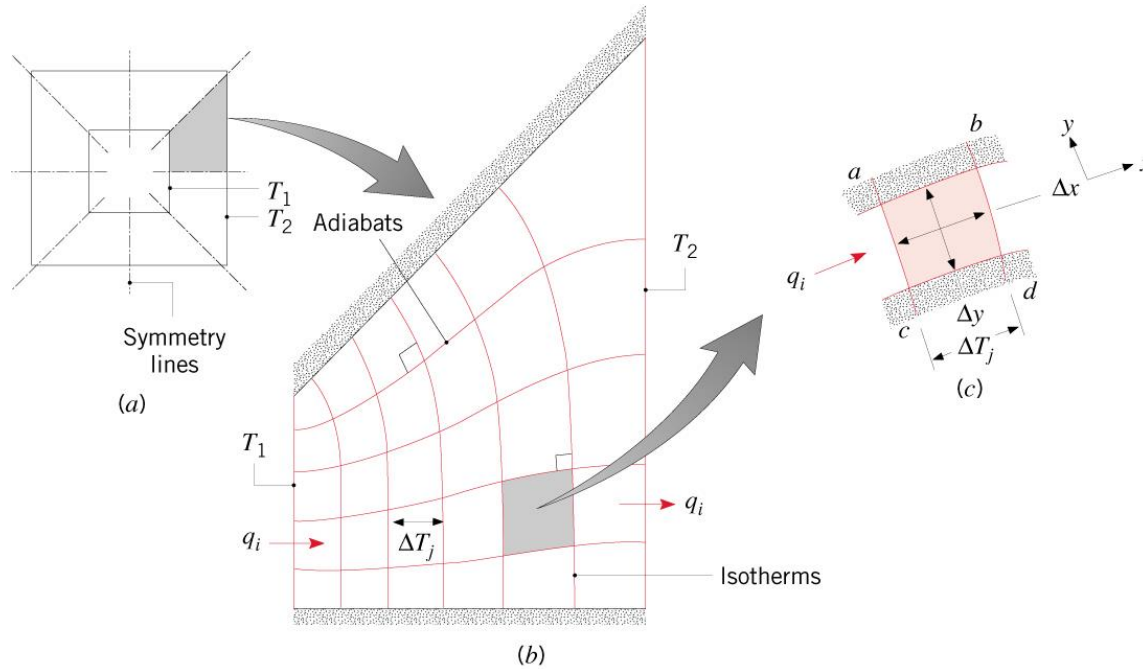
Métodos de Solução da Equação da Difusão de Calor

- Considerando regime permanente com propriedades constantes e com geração em coord. cartesianas
- Métodos de Solução:
 - Exata/Analítica: Método de Separação de Variáveis (possível transformada de Laplace ou Fourier
 - Limitado para problemas mais simples
 - Aproximado/Gráfico $\left(\begin{array}{c} \square \\ q = 0 \end{array} \right)$
 - Em desuso
 - Aproximado/Numérico: Volume-Diferenças Finitas ou Elementos Finitos
 - Mais adaptado para qualquer complexidade

Condições de Simetria para problemas gráficos e numéricos



Método Gráfico



– Determinação da taxa transferida de calor:

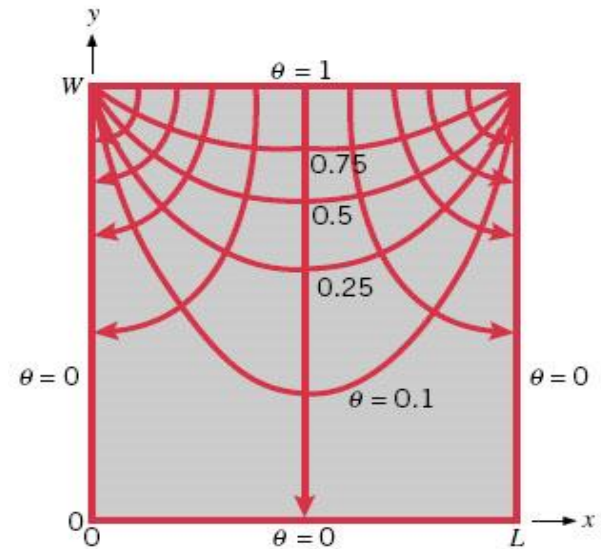
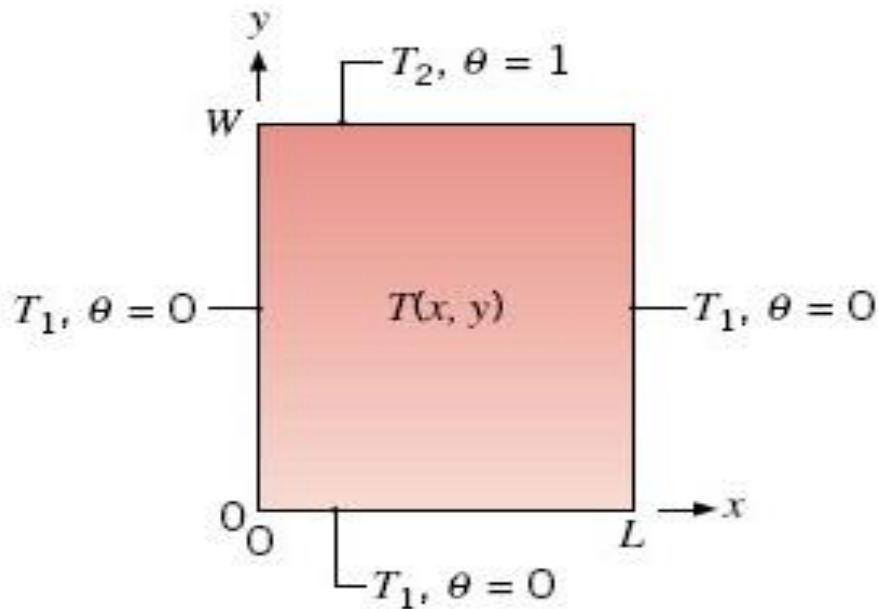
$$q \approx M q_i \approx M \left[k (\Delta y \cdot \ell) \frac{\Delta T_j}{\Delta x} \right] \approx \frac{M \ell}{N} k \Delta T_{1-2}$$

$$q' \approx \frac{M}{N} k \Delta T_{1-2}$$

M é o número de isofluxos e N o número de isoterms

Método Analítico

$$\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} = 0; \quad \theta = \frac{T - T_1}{T_2 - T_1}; \quad \frac{\partial^2 \theta}{\partial x^2} + \frac{\partial^2 \theta}{\partial y^2} = 0$$



Condições de contorno

$$\theta(0, y) = 0$$

$$\theta(L, y) = 0$$

$$\theta(x, 0) = 0$$

$$\theta(x, W) = 1$$

MÉTODO DE SEPARAÇÃO DE VARIÁVEIS

$$\theta(x, y) = X(x).Y(y)$$

Substituindo

$$X''.Y + Y''.X = 0 \quad \div XY$$

$$\frac{X''}{X} + \frac{Y''}{Y} = 0$$

$$-\frac{X''}{X} = \frac{Y''}{Y}$$

MÉTODO DE SEPARAÇÃO DE VARIÁVEIS

$$\theta(x, y) = X(x).Y(y)$$

Substituindo

$$X''.Y + Y''.X = 0 \quad \div XY$$

$$\frac{X''}{X} + \frac{Y''}{Y} = 0$$

$$-\frac{X''}{X} = \frac{Y''}{Y} = \lambda^2$$

Autovalor



$$-\frac{X''}{X} = \frac{Y''}{Y} = \lambda^2$$

Resultando

$$\begin{cases} Y'' - \lambda^2 Y = 0 & (I) \\ X'' + \lambda^2 X = 0 & (II) \end{cases}$$

De (I)

$$Y'' - \lambda^2 Y = 0$$

$$(D^2 - \lambda^2)Y = 0$$

$$(D - \lambda)(D + \lambda)Y = 0$$

$$Y(y) = c_1 e^{\lambda y} + c_2 e^{-\lambda y}$$

De (II) $X'' + \lambda^2 X = 0$

$$(D^2 + \lambda^2)X = 0 \quad ; m\sqrt{-\lambda^2} = \pm \lambda i$$

α_1 e α_2 são complexos

$$\alpha_1 = \lambda i \quad ; \quad \alpha_2 = -\lambda i$$

$$X(x) = d_1 e^{\alpha_1 x} + d_2 e^{\alpha_2 x} = d_1 e^{(a+bi)x} + d_2 e^{(a-bi)x}$$

$$X(x) = e^{ax} (d_1 e^{bix} + d_2 e^{-bix});$$

fórmula de Euler (Moivre): $e^{\pm ix} = \cos x \pm i \sin x$

$$X(x) = e^{ax} (d_1 (\cos bx + i \sin bx) + d_2 (\cos bx - i \sin bx))$$

$$X(x) = e^{ax} ((d_1 + d_2) \cos bx + i(d_1 - d_2) \sin bx)$$

$$X(x) = e^{ax} (c_3 \cos bx + c_4 \sin bx) \quad ; \text{como } a = 0$$

$$X(x) = c_3 \cos \lambda x + c_4 \sin \lambda x$$

Da solução Geral:

$$\theta(x, y) = X(x).Y(y)$$

$$\theta(x, y) = (c_3 \cos \lambda x + c_4 \text{sen } \lambda x)(c_1 e^{\lambda y} + c_2 e^{-\lambda y})$$

Aplicando as condições de contorno:

$$\theta(0, y) = 0 \rightarrow X(0) = c_3 \cos(\lambda \cdot 0) + c_4 \text{sen}(\lambda \cdot 0) \therefore c_3 = 0$$

$$\theta(x, 0) = 0 \rightarrow Y(0) = c_1 + c_2 = 0 \rightarrow c_1 = -c_2$$

$$\theta(L, y) = 0 \rightarrow X(L) = 0 = c_4 \text{sen}(\lambda \cdot L); \text{ como } c_4 \neq 0$$
$$\text{sen}(\lambda \cdot L) = 0$$

$$\text{então } \lambda \cdot L = n\pi; n = 1, 2, 3, \dots, \infty$$

$$\lambda_n = \frac{n\pi}{L} \text{ (autovalores)}$$

$$c_3 = 0$$

$$c_1 = -c_2$$

$$\lambda_n = \frac{n\pi}{L} \text{ (autovalores)}$$

Da solução Geral:

$$\theta(x, y) = X(x).Y(y)$$

$$\theta_n(x, y) = (c_1 c_4)_n (\text{sen } \lambda_n x) (e^{\lambda_n y} - e^{-\lambda_n y})$$

$$\theta_n(x, y) = c'_n (\text{sen } \lambda_n x) (e^{\lambda_n y} - e^{-\lambda_n y})$$

como $(e^{\lambda_n y} - e^{-\lambda_n y}) = 2 \text{senh}(\lambda_n y)$

$$\theta_n(x, y) = c_n (\text{sen } \lambda_n x) (\text{senh } \lambda_n y)$$

Tem-se um infinito número de soluções que satisfazem a equação diferencial e as condições de contorno sendo cada solução linearmente independente. Como o problema é linear a solução geral será a superposição de todas estas soluções na forma:

$$\theta(x, y) = \sum_{n=1}^{\infty} c_n (\text{sen } \lambda_n x)(\text{senh } \lambda_n y)$$

Para determinar-se a constante c_n , aplica-se a última condição de contorno,

$$\theta(x, W) = 1$$

$$\theta(x, W) = 1 = \sum_{n=1}^{\infty} c_n (\text{sen } \lambda_n x)(\text{senh } \lambda_n W)$$

Utilizando as propriedades de funções ortogonais:

$$\int_a^b g_m(x)g_n(x)dx = 0 \quad p / m \neq n$$

Multiplicando de ambos os lados da equação da pag. anterior

$$\int_0^L 1 \cdot \text{sen } \lambda_m x dx = \int_0^L \sum_{n=1}^{\infty} c_n (\text{sen } \lambda_n x) (\text{senh } \lambda_n W) \text{sen } \lambda_m x dx$$

$$c_n = \frac{\int_0^L 1 \cdot \text{sen } \lambda_m x dx}{\int_0^L \sum_{n=1}^{\infty} (\text{sen } \lambda_n x)^2 (\text{senh } \lambda_n W) dx} = \frac{\int_0^L 1 \cdot \text{sen } \lambda_m x dx}{(\text{senh } \lambda_n W) \int_0^L \sum_{n=1}^{\infty} (\text{sen } \lambda_n x)^2 dx}$$

$$c_n = \frac{2[(-1)^{n+1} + 1]}{(\text{senh } \lambda_n W) n \pi}, \quad n = 1, 2, \dots, \infty$$

Solução geral

$$\theta(x, y) = \frac{2}{\pi} \sum_{n=1}^{\infty} (-1)^{n+1} (\text{sen } \lambda_n x) \frac{(\text{senh } \lambda_n y)}{(\text{senh } \lambda_n W)}$$

Métodos de Solução da Equação da Difusão de Calor

- Considerando regime permanente com propriedades constantes e com geração em coord. cartesianas

$$\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\dot{q}(x, y)}{k} = 0$$

- **Métodos de Solução:**

- **Exata/Analítica:** Método de Separação de Variáveis (possível transformada de Laplace ou Fourier)

- Limitado para problemas mais simples

- **Approximado/Gráfico** $\left(\dot{q} = 0 \right)$

- Em desuso

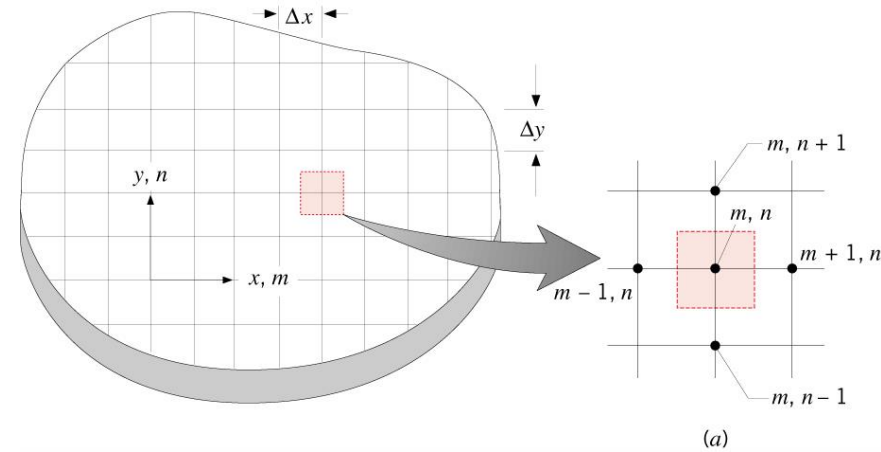
- **Approximado/Numérico:** Volume-Diferenças Finitas ou Elementos Finitos

- Mais adaptado para qualquer complexidade

– Métodos Numérico:

– Diferenças Finitas, Método dos Balanços ou Volumes Finitos

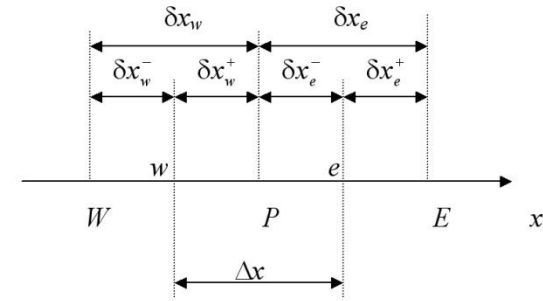
– Elementos Finitos



Método dos Volumes Finitos

$$\int_V \left(\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) \right) dV = 0$$

$$\int_t^{t+\Delta t} dt \int_w^e dx \int_s^n dy$$

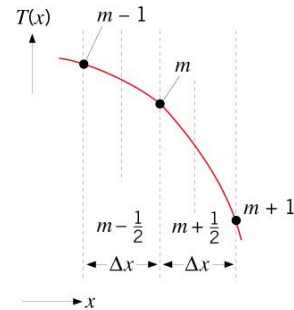


$$\int_w^e \int_s^n \rho u_s \frac{\partial \phi}{\partial s} dx dy = \int_w^e \int_s^n \frac{\partial}{\partial x} \left(\Gamma \frac{\partial \phi}{\partial x} \right) dx dy + \int_w^e \int_s^n \frac{\partial}{\partial y} \left(\Gamma \frac{\partial \phi}{\partial y} \right) dx dy + \int_w^e \int_s^n S dx dy$$

Método de Diferenças Finitas

$$\left. \frac{\partial T}{\partial x} \right|_{m-1/2, n} = \frac{T_{m, n} - T_{m-1, n}}{\Delta x}$$

$$\left. \frac{\partial T}{\partial x} \right|_{m+1/2, n} = \frac{T_{m+1, n} - T_{m, n}}{\Delta x}$$



(b)

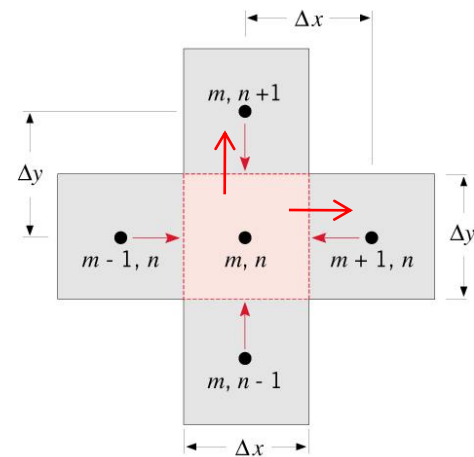
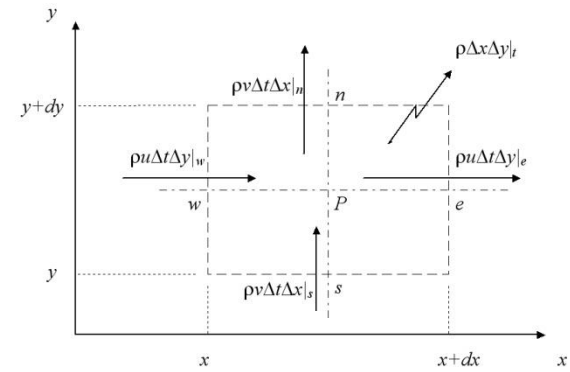
Método do Balanço de Energia

$$\dot{E}_{in} + \dot{E}_g = 0$$

Em Regime Permanente

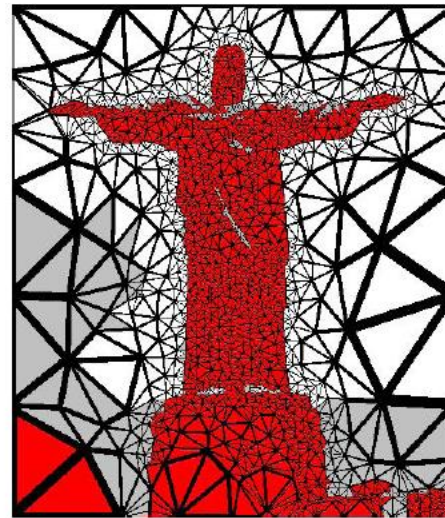
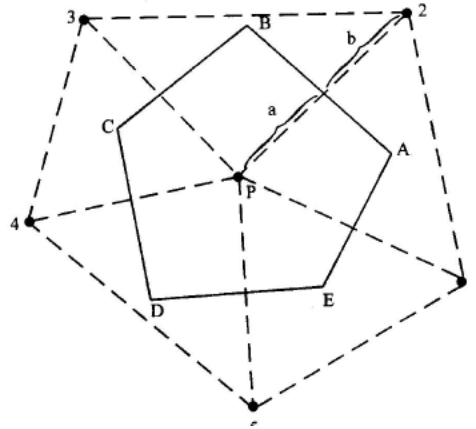
$$\sum_{i=1}^4 q_{(i) \rightarrow (m,n)} + \dot{q}(\Delta x \cdot \Delta y \cdot \ell) = 0$$

$$q_{(m-1,n) \rightarrow (m,n)} = k(\Delta y \cdot \ell) \frac{T_{m-1,n} - T_{m,n}}{\Delta x}$$



(4.31)

Diagrama de Voronoi



Malhas

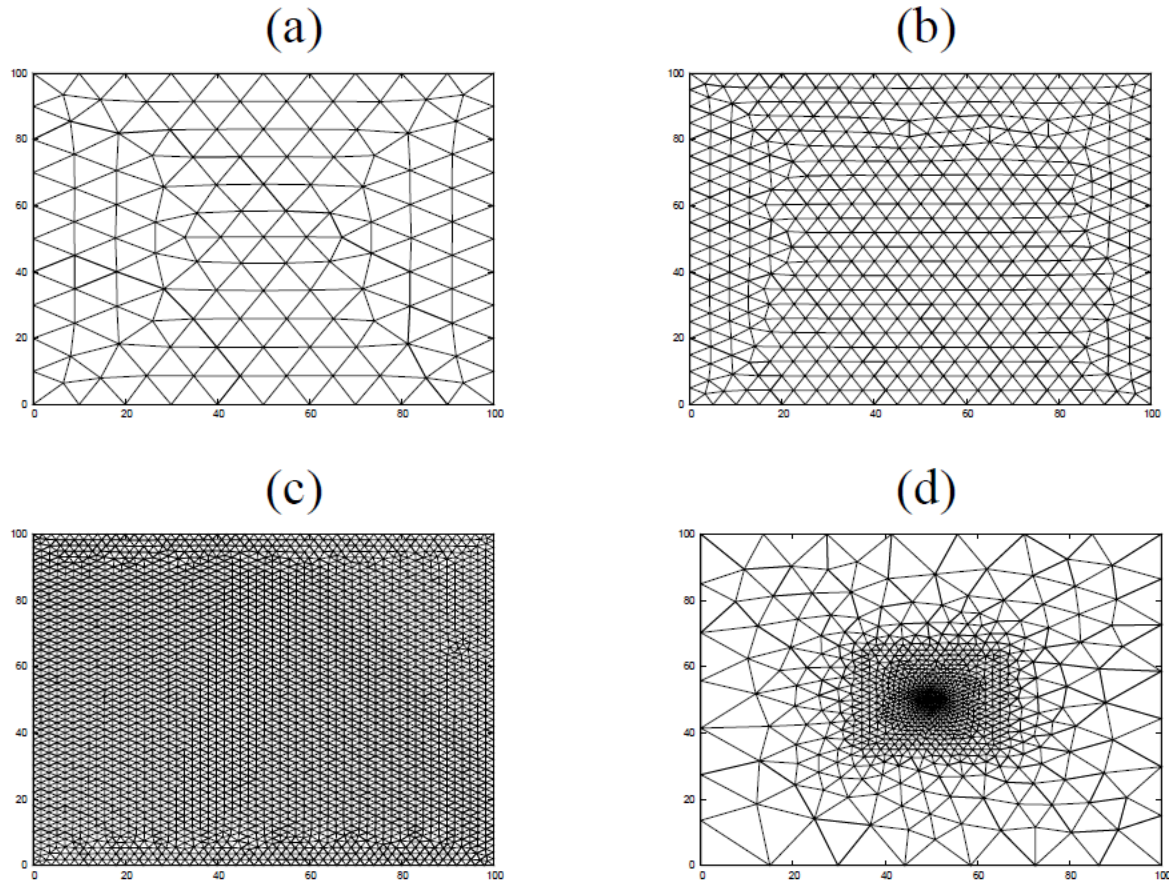
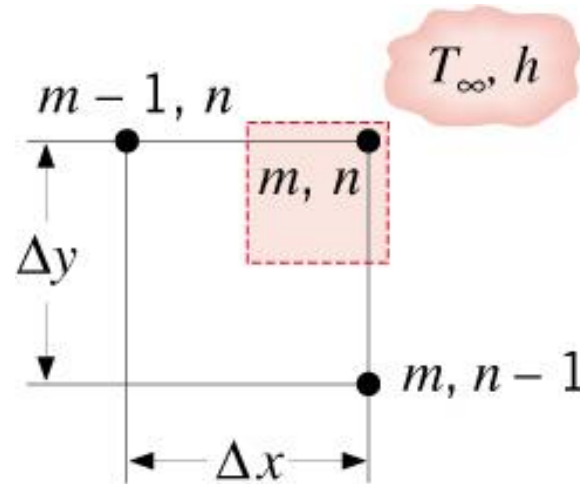


Figura 14: Malhas para a simulação de Volumes Finitos

- A summary of finite-difference equations for common nodal regions is provided in Table 4.2. Consider an *external corner with convection heat transfer*.



$$q_{(m-1,n) \rightarrow (m,n)} + q_{(m,n-1) \rightarrow (m,n)} + q_{(\infty) \rightarrow (m,n)} = 0$$

$$k \left(\frac{\Delta y}{2} \cdot \ell \right) \frac{T_{m-1,n} - T_{m,n}}{\Delta x} + k \left(\frac{\Delta x}{2} \cdot \ell \right) \frac{T_{m,n-1} - T_{m,n}}{\Delta y} + h \left(\frac{\Delta x}{2} \cdot \ell \right) (T_\infty - T_{m,n}) + h \left(\frac{\Delta y}{2} \cdot \ell \right) (T_\infty - T_{m,n}) = 0$$

or, with $\Delta x = \Delta y$,

$$T_{m-1,n} + T_{m,n-1} + 2 \frac{h\Delta x}{k} T_\infty - 2 \left(\frac{h\Delta x}{k} + 1 \right) T_{m,n} = 0$$

(4.43)

Métodos de Solução

- **Inversão da Matriz:** Expression of system of N finite-difference equations for N unknown nodal temperatures as:

$$[A][T] = [C] \quad (4.48)$$

Coefficient Matrix (NxN)
Solution Vector (T₁, T₂, ... T_N)
Right-hand Side Vector of Constants (C₁, C₂... C_N)

Solution →

$$[T] = [A]^{-1} [C] \quad (4.49)$$

Inverse of Coefficient Matrix

- **Gauss-Seidel** : Each finite-difference equation is written in **explicit form**, such that its unknown nodal temperature appears alone on the left-hand side:

$$T_i^{(k)} = \frac{C_i}{a_{ii}} - \sum_{j=1}^{i-1} \frac{a_{ij}}{a_{ii}} T_j^{(k)} - \sum_{j=i+1}^N \frac{a_{ij}}{a_{ii}} T_j^{(k-1)} \quad (4.51)$$

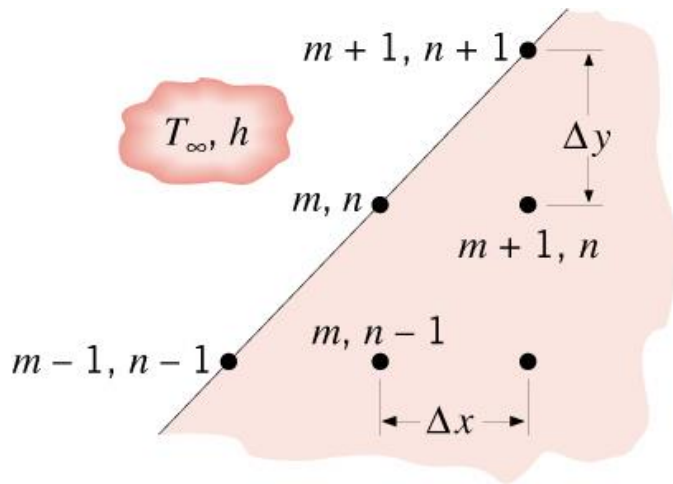
where $i = 1, 2, \dots, N$ and k is the level of iteration.

Iteration proceeds until satisfactory convergence is achieved for all nodes:

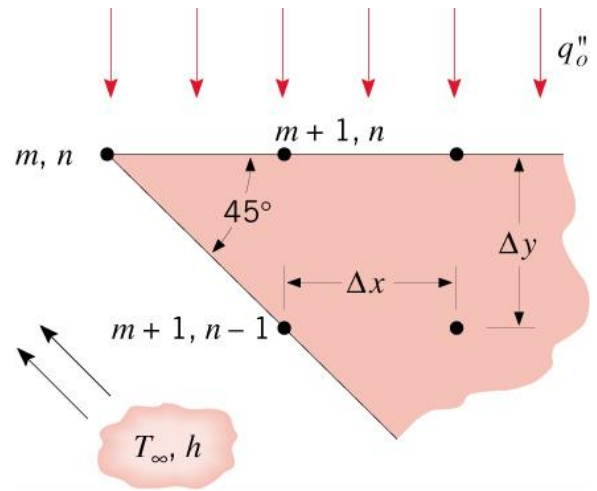
$$\left| T_i^{(k)} - T_i^{(k-1)} \right| \leq \varepsilon$$

- What measures may be taken to insure that the results of a finite-difference solution provide an accurate prediction of the temperature field?

Problem 4.39: Finite-difference equations for (a) nodal point on a diagonal surface and (b) tip of a cutting tool.

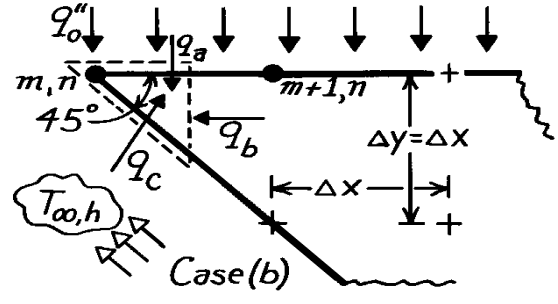
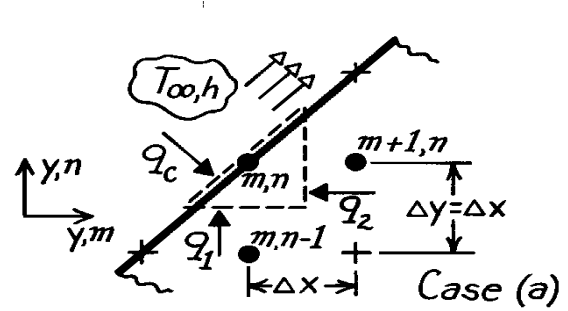


(a) Diagonal surface



(b) Cutting tool.

Schematic:



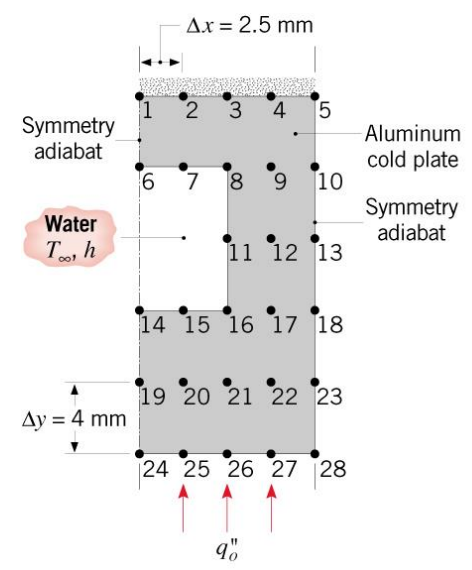
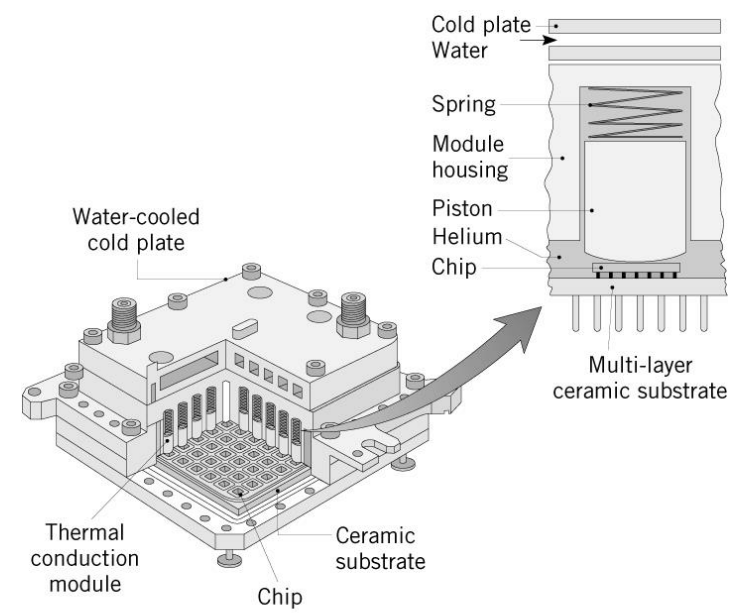
ASSUMPTIONS: (1) Steady-state, 2-D conduction, (2) Constant properties

Problem 4.76: Analysis of cold plate used to thermally control IBM multi-chip, thermal conduction module.

Features:

- Heat dissipated in the chips is transferred by conduction through spring-loaded aluminum pistons to an aluminum cold plate.
- Nominal operating conditions may be assumed to provide a uniformly distributed heat flux of at the base of the cold plate.
- Heat is transferred from the cold plate by water flowing through channels in the cold plate.

Find: (a) Cold plate temperature distribution for the prescribed conditions. (b) Options for operating at larger power levels while remaining within a maximum cold plate temperature of 40°C.



ANALYSIS: Finite-difference equations must be obtained for each of the 28 nodes. Applying the energy balance method to regions 1 and 5, which are similar, it follows that

$$\text{Node 1: } (\Delta y/\Delta x)T_2 + (\Delta x/\Delta y)T_6 - [(\Delta y/\Delta x) + (\Delta x/\Delta y)]T_1 = 0$$

$$\text{Node 5: } (\Delta y/\Delta x)T_4 + (\Delta x/\Delta y)T_{10} - [(\Delta y/\Delta x) + (\Delta x/\Delta y)]T_5 = 0$$

Nodal regions 2, 3 and 4 are similar, and the energy balance method yields a finite-difference equation of the form

Nodes 2,3,4:

$$(\Delta y/\Delta x)(T_{m-1,n} + T_{m+1,n}) + 2(\Delta x/\Delta y)T_{m,n-1} - 2[(\Delta y/\Delta x) + (\Delta x/\Delta y)]T_{m,n} = 0$$

Energy balances applied to the remaining combinations of similar nodes yield the following finite-difference equations.

$$\begin{aligned} \text{Nodes 6, 14: } (\Delta x/\Delta y)T_1 + (\Delta y/\Delta x)T_7 - [(\Delta x/\Delta y) + (\Delta y/\Delta x) + (h\Delta x/k)]T_6 &= -(h\Delta x/k)T_\infty \\ (\Delta x/\Delta y)T_{19} + (\Delta y/\Delta x)T_{15} - [(\Delta x/\Delta y) + (\Delta y/\Delta x) + (h\Delta x/k)]T_{14} &= -(h\Delta x/k)T_\infty \end{aligned}$$

$$\begin{aligned} \text{Nodes 7, 15: } (\Delta y/\Delta x)(T_6 + T_8) + 2(\Delta x/\Delta y)T_2 - 2[(\Delta y/\Delta x) + (\Delta x/\Delta y) + (h\Delta x/k)]T_7 &= -(2h\Delta x/k)T_\infty \\ (\Delta y/\Delta x)(T_{14} + T_{16}) + 2(\Delta x/\Delta y)T_{20} - 2[(\Delta y/\Delta x) + (\Delta x/\Delta y) + (h\Delta x/k)]T_{15} &= -(2h\Delta x/k)T_\infty \end{aligned}$$

$$\begin{aligned} \text{Nodes 8, 16: } & (\Delta y/\Delta x)T_7 + 2(\Delta y/\Delta x)T_9 + (\Delta x/\Delta y)T_{11} + 2(\Delta x/\Delta y)T_3 - [3(\Delta y/\Delta x) + 3(\Delta x/\Delta y) \\ & + (h/k)(\Delta x + \Delta y)]T_8 = -(h/k)(\Delta x + \Delta y)T_\infty \\ & (\Delta y/\Delta x)T_{15} + 2(\Delta y/\Delta x)T_{17} + (\Delta x/\Delta y)T_{11} + 2(\Delta x/\Delta y)T_{21} - [3(\Delta y/\Delta x) + 3(\Delta x/\Delta y) \\ & + (h/k)(\Delta x + \Delta y)]T_{16} = -(h/k)(\Delta x + \Delta y)T_\infty \end{aligned}$$

$$\text{Node 11: } (\Delta x/\Delta y)T_8 + (\Delta x/\Delta y)T_{16} + 2(\Delta y/\Delta x)T_{12} - 2[(\Delta x/\Delta y) + (\Delta y/\Delta x) + (h\Delta y/k)]T_{11} = -(2h\Delta y/k)T_\infty$$

Nodes 9, 12, 17, 20, 21, 22:

$$(\Delta y/\Delta x)T_{m-1,n} + (\Delta y/\Delta x)T_{m+1,n} + (\Delta x/\Delta y)T_{m,n+1} + (\Delta x/\Delta y)T_{m,n-1} - 2[(\Delta x/\Delta y) + (\Delta y/\Delta x)]T_{m,n} = 0$$

Nodes 10, 13, 18, 23:

$$(\Delta x/\Delta y)T_{n+1,m} + (\Delta x/\Delta y)T_{n-1,m} + 2(\Delta y/\Delta x)T_{m-1,n} - 2[(\Delta x/\Delta y) + (\Delta y/\Delta x)]T_{m,n} = 0$$

$$\text{Node 19: } (\Delta x/\Delta y)T_{14} + (\Delta x/\Delta y)T_{24} + 2(\Delta y/\Delta x)T_{20} - 2[(\Delta x/\Delta y) + (\Delta y/\Delta x)]T_{19} = 0$$

$$\begin{aligned} \text{Nodes 24, 28: } & (\Delta x/\Delta y)T_{19} + (\Delta y/\Delta x)T_{25} - [(\Delta x/\Delta y) + (\Delta y/\Delta x)]T_{24} = -(q_o''\Delta x/k) \\ & (\Delta x/\Delta y)T_{23} + (\Delta y/\Delta x)T_{27} - [(\Delta x/\Delta y) + (\Delta y/\Delta x)]T_{28} = -(q_o''\Delta x/k) \end{aligned}$$

Nodes 25, 26, 27:

$$(\Delta y/\Delta x)T_{m-1,n} + (\Delta y/\Delta x)T_{m+1,n} + 2(\Delta x/\Delta y)T_{m,n+1} - 2[(\Delta x/\Delta y) + (\Delta y/\Delta x)]T_{m,n} = -(2q_o''\Delta x/k)$$

Evaluating the coefficients and solving the equations simultaneously, the steady-state temperature distribution (°C), tabulated according to the node locations, is:

23.77	23.91	24.27	24.61	24.74
23.41	23.62	24.31	24.89	25.07
		25.70	26.18	26.33
28.90	28.76	28.26	28.32	28.35
30.72	30.67	30.57	30.53	30.52
32.77	32.74	32.69	32.66	32.65

(b) For the prescribed conditions, the maximum allowable temperature ($T_{24} = 40^\circ\text{C}$) is reached when

$$q_0'' = 1.407 \times 10^5 \text{ W/m}^2 \text{ (14.07 W/cm}^2\text{)}.$$

Options for extending this limit could include use of a copper cold plate ($k \approx 400 \text{ W/m}\cdot\text{K}$) and/or increasing the convection coefficient associated with the coolant.

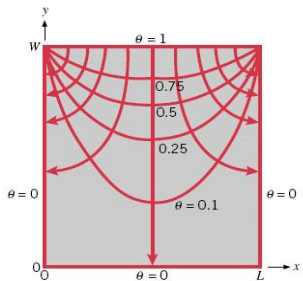
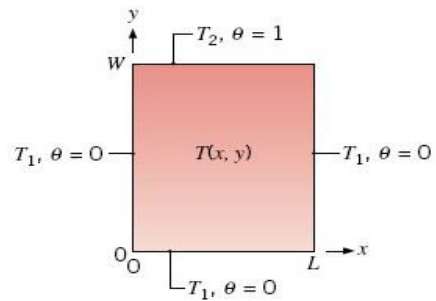
With $k = 400 \text{ W/m}\cdot\text{K}$, a value of $q_0'' = 17.37 \text{ W/cm}^2$ may be maintained.

With $k = 400 \text{ W/m}\cdot\text{K}$ and $h = 10,000 \text{ W/m}^2\cdot\text{K}$ (a practical upper limit), $q_0'' = 28.65 \text{ W/cm}^2$.

Additional, albeit small, improvements may be realized by relocating the coolant channels closer to the base of the cold plate.

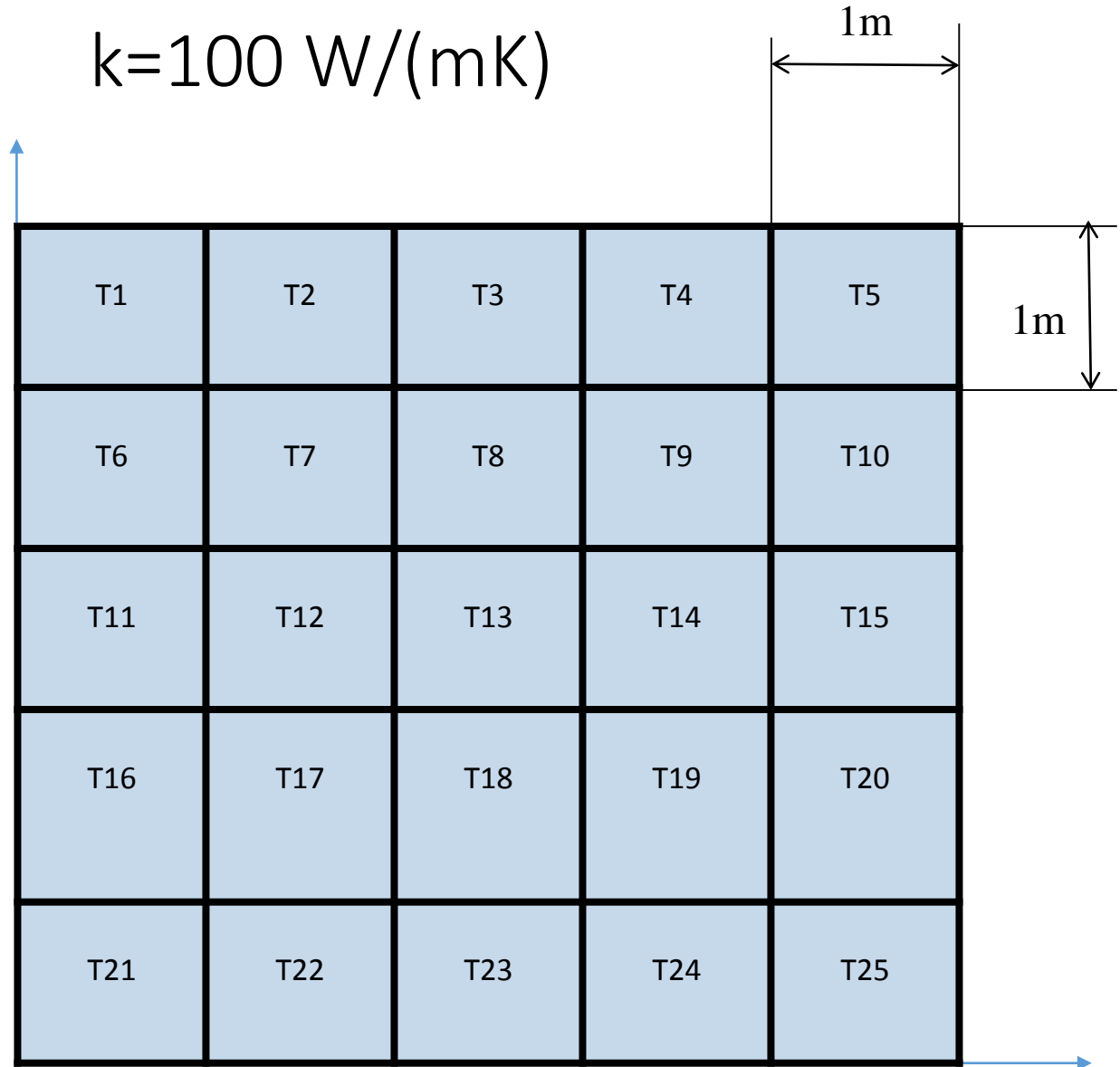
Exercício 1

$$k=100 \text{ W/(mK)}$$



Calcular:

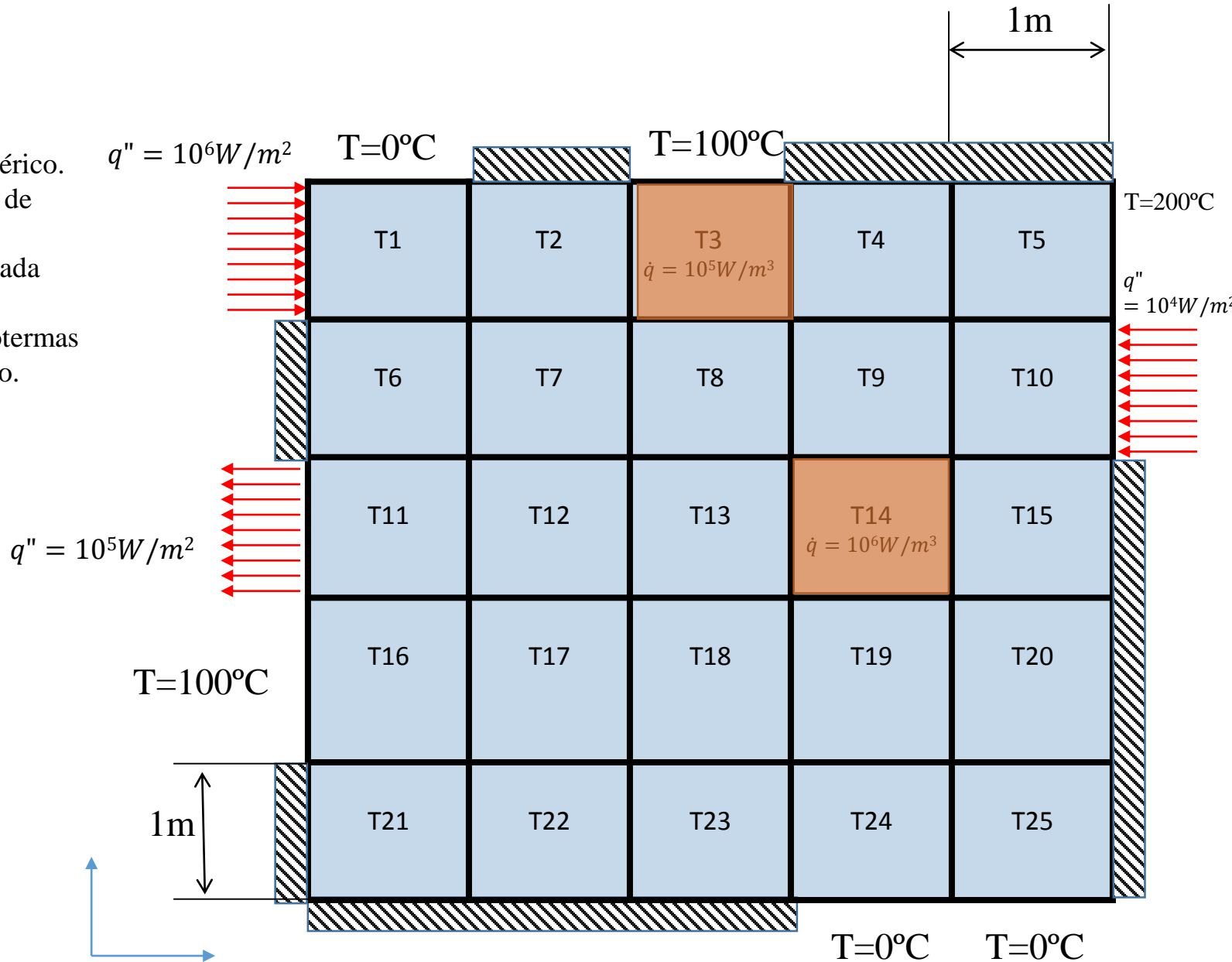
- 1) T_1, \dots, T_{25} . Analítico e numérico.
- 2) Comparar resultados
- 3) Calcular fluxos de calor nas paredes e no centro.
- 4) Desenhar as isothermas e linhas de fluxo.



Exercício 2 $k=100 \text{ W}/(\text{mK}), w=1\text{m}$

Calcular:

- 1) T_1, \dots, T_6 . Numérico.
- 2) Calcular fluxos de calor vertical e horizontal em cada volume.
- 3) Desenhar as isotermas e linhas de fluxo.



Exercício 3 $k=100 \text{ W/(mK)}$, $w=1\text{m}$ NÃO COBRADO

Calcular:

- 1) T_1, \dots, T_6 . Numérico.
- 2) Calcular fluxos de calor vertical e horizontal em cada volume.
- 3) Desenhar as isotermas e linhas de fluxo.

