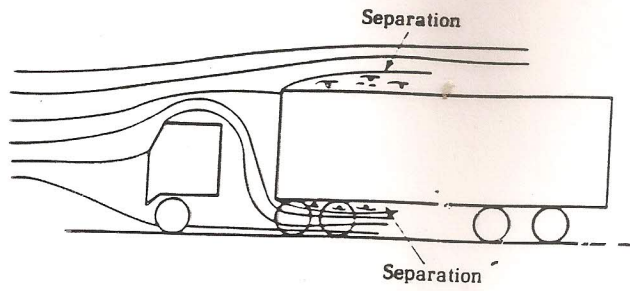
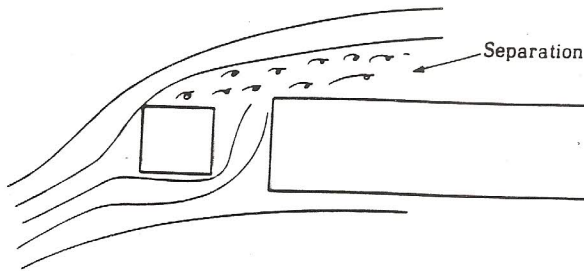


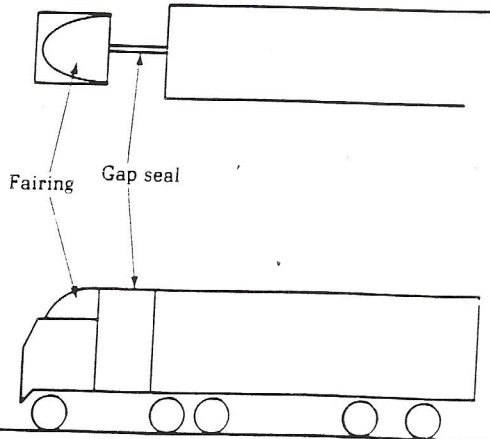
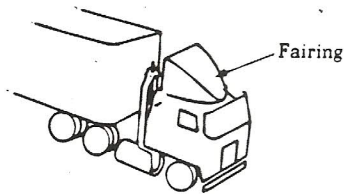
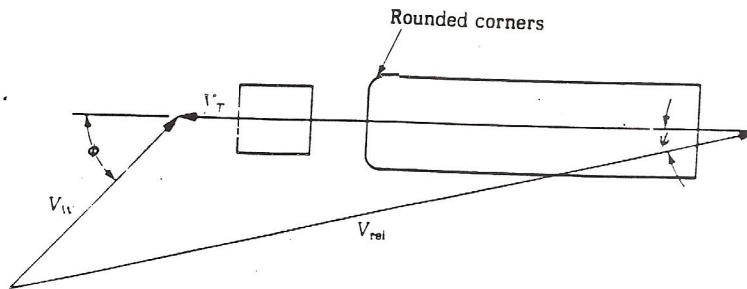
**FIGURE 6.22**  
Plan and profile views of streamlines about a cab-over-engine tractor-trailer combination.



**FIGURE 6.23**  
Plan view of the flow pattern on a yawed truck.



**FIGURE 6.24**  
Wind velocity diagram for analysis of flow about a truck.



**FIGURE 6.25**  
A fairing and gap seal attached to a tractor-trailer truck.

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as the sphere. In this respect is tentatively:

$$\frac{R_{crit}}{R_{crit0}} = \sqrt{\frac{V^2}{V^2 + (k \cdot u)^2}} = \frac{1}{\sqrt{1 + (k \cdot u/V)^2}}$$

where  $k < 1$  represents a suitable integration constant. This effect is demonstrated in figure 24 at (t); the critical Reynolds number of the sphere as tested is shifted from  $R_{d,crit} = 2.9 \cdot 10^5$  (without rotation) to  $2 \cdot 10^5$  as shown in the graph at  $u/V \approx 1.8$ ; hence  $k \approx 0.6$ . The sphere also exhibits the centrifugal effect at "c". Note that the low sphere value  $C_{D0} = 0.38$  in subcritical condition (instead of  $\approx 0.47$ ) is caused by type and size of the support in these tests by means of a rotating shaft.

*The Human Body* is similar in aerodynamic shape to a cylinder with a length ratio  $h/d$  between 4 and 7. Since human beings vary very much in size and proportions, selection of a reference area is difficult. Figure 25, therefore, presents the drag of an average man in the form of drag area  $D/q$ . The drag is predominantly a function of the projected frontal area in the various positions tested. Based on estimated areas, drag coefficients can be determined for the standing positions between  $C_{D0} = 1.0$  and 1.3. Without clothing, the drag is between 5 and 10% less than listed.

*How Fast a Man Falls.* After bailing out of an airplane, and before releasing the parachute, the body of a man accelerates to a terminal velocity the magnitude of which can be derived from  $W = D = q(D/q)$ . Near sea level (where  $\rho = 0.0024 \text{ lb sec}^2/\text{ft}^2$ ), the falling speed of a man with  $W = 180 \text{ lb}$ , is accordingly  $V_{ft/sec} \approx 400 \sqrt{(D/q)}$ . Employing the drag areas as listed in figure 25 (between 1.2 and 9.0  $\text{ft}^2$ ), speeds between 130 and 370  $\text{ft/sec}$  are thus obtained. Terminal velocities are reported (32,c) between 150 and 180  $\text{ft/sec}$  "near sea level"; without specification as to position and attitude during free fall. Another source (32.e) gives a drag area of 5  $\text{ft}^2$  for a "rolling and somersaulting" man. To give a certain scale to all these numbers, it is mentioned that the drag area of a typical fighter airplane is in the order of 6  $\text{ft}^2$ .

*The Drag of Ski-Runners* has been tested in wind tunnels. In upright position (going down a slope) a drag area  $D/q = 5.5 \text{ ft}^2$  is found (32,a) in a smooth wooden model. A similar value ( $\approx 6.5 \text{ ft}^2$ ) can be derived from (32,b) on the basis of an estimated frontal area in the order of 7  $\text{ft}^2$ . Both sources also give results on drag and lift of a ski-jumping man. In the typical "flying" position, with the body leaning forward against and onto the air, the lift area (including the contribution of the skis) is in the order of  $L/q = 2.5 \text{ ft}^2$ ; the maximum lift/drag ratio is in the order of "1".

### 6. DRAG OF VARIOUS TYPES OF PLATES

All that is said in the preceding section about the critical effect of the boundary layer upon the drag of spheres, applies in principle to all sufficiently rounded bodies, such as the strut sections for instance in Chapter VI. On the other hand, bodies with sharp edges, such as disks and plates in a flow normal to their surfaces, do not show any critical drag decrease. The pressure gradient around the sharp edges would necessarily be extremely high for a flow pattern attached to the rear of a plate — that is, theoretically from  $\Delta p/q = -\infty$  at the edge to  $+1$  at the rear stagnation point. No boundary layer, whether laminar or turbulent, can follow the way around the edges of such plates.

*Small R Numbers.* Figure 26 shows the drag coefficient of disks and square plates in normal flow, as a function of Reynolds number. Below  $R_d = 100$ , there is the regime of predominantly viscous flow as discussed in the beginning of this chapter. Approximately at  $R_d = 300$ , the drag coefficient of the disk shows a peak, as reported from two independent sources. Observation of the flow pattern (35) proves this peak to be due to a change in the pattern of the vortex system behind the body.

*Turbulence Effect.* Above  $R_d = 1000$ , the drag coefficient of disks (and other plates) is practically constant up to the highest Reynolds numbers ever tested (approaching  $10^7$ ). Because of this stability,

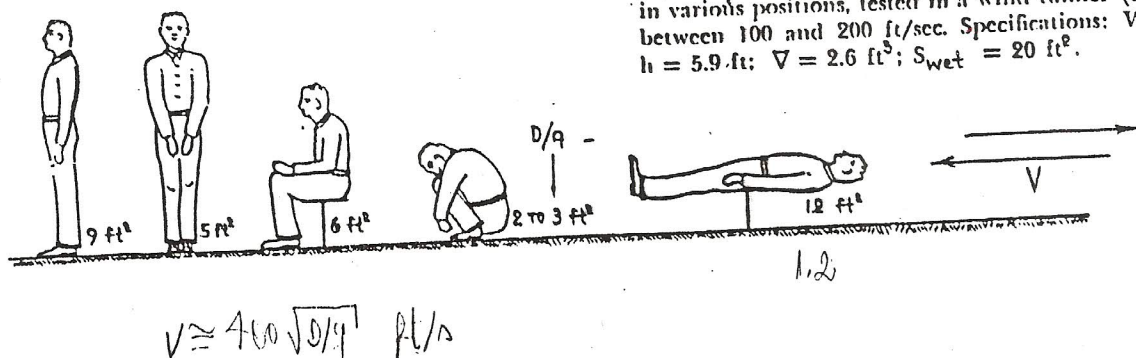


Figure 25. Drag areas ( $D/q$  in  $\text{ft}^2$ ) of an average man in various positions, tested in a wind tunnel (31) at speeds between 100 and 200  $\text{ft/sec}$ . Specifications:  $W = 165 \text{ lb}$ ;  $h = 5.9 \text{ ft}$ ;  $\nabla = 2.6 \text{ ft}^3$ ;  $S_{wet} = 20 \text{ ft}^2$ .